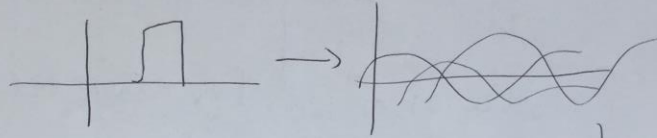




## Ch 7 Super-position of Quasi-Parallel Plane Waves (1)

- Any Field can be described as a superposition of many plane wave fields. (pulses).



### 7.1 Intensity of superimposed Plane Waves

- Arbitrary waveforms =  $\sum$  plane waves with different  $\vec{k}$ , amplitudes, phases, frequencies, polarizations.

$$\vec{E}(\vec{r}, t) = \sum_j \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)}$$

$$\vec{B}(\vec{r}, t) = \sum_j \vec{B}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)}$$

$$= \sum_j \frac{\vec{k}_j \times \vec{E}_j}{\omega_j} e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)}$$

$$\vec{S}(\vec{r}, t) = \text{Re} \{ \vec{E}(\vec{r}, t) \} \times \text{Re} \{ \vec{B}(\vec{r}, t) \} \quad (2)$$

$$= \sum_{j,m} \frac{1}{\omega_m \mu_0} \text{Re} \left\{ \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \right\} \times \text{Re} \left\{ \vec{k}_m \times \vec{E}_m e^{i(\vec{k}_m \cdot \vec{r} - \omega_m t)} \right\} \quad (7.3)$$

(cross) ←

[Notice  $j, m$  for two series terms].

Simplifying Assumptions to begin with:

- A1 1) Time-Average  $\vec{S}$  to remove fluctuations that change on scale of optical frequencies. ?? Why
- A2 2) Assume all plane waves travel  $\approx$  parallel to each other.
- A3 3) let  $\vec{k}_m$  be real.

Use  $\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B})$  (3)

on (7.3)

$$\vec{A} = \text{Re} \left\{ \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \right\}$$

$$\vec{B} = \vec{k}_m \quad ; \quad \vec{C} = \text{Re} \left\{ \vec{E}_m e^{i(\vec{k}_m \cdot \vec{r} - \omega_m t)} \right\}$$

$$\therefore \vec{S}(\vec{r}, t) = \sum_{j,m} \frac{1}{\omega_m \mu_0} \left[ \vec{k}_m (\text{Re} \left\{ \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \right\} \cdot \right.$$

$$\left. \text{Re} \left\{ \vec{E}_m e^{i(\vec{k}_m \cdot \vec{r} - \omega_m t)} \right\} \right)$$

$$- \text{Re} \left\{ \vec{E}_m e^{i(\vec{k}_m \cdot \vec{r} - \omega_m t)} \right\} (\text{Re} \left\{ \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \right\} \cdot \vec{k}_m) \right]$$

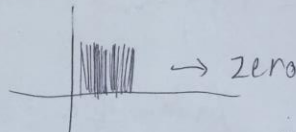
Using A2: Last term  $\rightarrow 0$  because

$$\vec{k}_m \perp \vec{E}_j$$

$$\vec{S}(\vec{r}, t) = \sum_{j,m} \frac{k_m}{\omega_m \mu_0} \left\{ \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} + \vec{E}_j^* e^{-i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \right\} \cdot \left\{ \frac{\vec{E}_m e^{i(\vec{k}_m \cdot \vec{r} - \omega_m t)} + \vec{E}_m^* e^{-i(\vec{k}_m \cdot \vec{r} - \omega_m t)}}{2} \right\}$$

$$\vec{S}(\vec{r}, t) = \sum_{j,m} \frac{k_m}{4W_m \mu_0} \left\{ \begin{aligned} & \vec{E}_j \cdot \vec{E}_m e^{i[(\vec{k}_j + \vec{k}_m) \cdot \vec{r} - (\omega_j + \omega_m)t]} \\ & + \vec{E}_j \cdot \vec{E}_m e^{-i[(\vec{k}_j + \vec{k}_m) \cdot \vec{r} - (\omega_j + \omega_m)t]} \\ & + \vec{E}_j \cdot \vec{E}_m e^{i[(\vec{k}_j - \vec{k}_m) \cdot \vec{r} - (\omega_j - \omega_m)t]} \\ & + \vec{E}_j \cdot \vec{E}_m e^{-i[(\vec{k}_j - \vec{k}_m) \cdot \vec{r} - (\omega_j - \omega_m)t]} \end{aligned} \right\} \quad (4)$$

Using A1: all  $(\omega_j + \omega_m)t$  terms oscillate quickly and time-average to zero.



We keep  $(\omega_j - \omega_m)t$ . It oscillates slowly. Why??

Assume  $n \approx \text{constant}$ .

$$\therefore k_m / W_m \mu_0 \approx n E_0 c$$

$$\langle \vec{S}(\vec{r}, t) \rangle_{osc} = \frac{n \epsilon_0 c}{2} \sum_{j, m} \frac{\vec{E}_j \cdot \vec{E}_m^* e^{i[(\vec{k}_j - \vec{k}_m) \cdot \vec{r} - (\omega_j - \omega_m)t]} + \vec{E}_j^* \cdot \vec{E}_m e^{-i[(\vec{k}_j - \vec{k}_m) \cdot \vec{r} - (\omega_j - \omega_m)t]}}{2} \quad (5)$$

$$= \frac{n \epsilon_0 c}{2} \text{Re} \left\{ \sum_j \vec{E}_j e^{i(\vec{k}_j \cdot \vec{r} - \omega_j t)} \sum_m \vec{E}_m^* e^{-i(\vec{k}_m \cdot \vec{r} - \omega_m t)} \right\}$$

$$= \frac{n \epsilon_0 c}{2} \text{Re} \left\{ \vec{E}(\vec{r}, t) \cdot \vec{E}^*(\vec{r}, t) \right\}$$

$$\therefore \underline{I}(\vec{r}, t) = \frac{n \epsilon_0 c}{2} \vec{E}(\vec{r}, t) \cdot \vec{E}^*(\vec{r}, t)$$

Parallel or nearly parallel  
 $\vec{k}_j$

## 7.2 Group vs. Phase Velocity: (6)

### Sum of two plane waves

$$\vec{E}_1 = \vec{E}_0 e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} \quad \vec{E}_2 = \vec{E}_0 e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)}$$

Equal Amplitudes.

phase velocities:  $V_1 = \frac{\omega_1}{k_1}$  ;  $V_2 = \frac{\omega_2}{k_2}$

$$\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} + \vec{E}_0 e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)}$$



interference occurs.  
constructive + destructive.

$$\begin{aligned} I(\vec{r}, t) &= n \epsilon_0 c \vec{E}(\vec{r}, t) \cdot \vec{E}^*(\vec{r}, t) \\ &= \frac{n \epsilon_0 c}{2} \left[ \vec{E}_0 e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} + \vec{E}_0 e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)} \right] \\ &\quad \cdot \left[ \vec{E}_0^* e^{-i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} + \vec{E}_0^* e^{-i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)} \right] \end{aligned}$$



$$= \frac{n \epsilon_0 c}{2} \left[ \vec{E}_0 \cdot \vec{E}_0^* + \vec{E}_0 \cdot \vec{E}_0^* e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} - i(\vec{k}_2 \cdot \vec{r} - \omega_2 t) \right. \\ \left. + \vec{E}_0 \cdot \vec{E}_0^* e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)} - i(\vec{k}_1 \cdot \vec{r} - \omega_1 t) \right. \\ \left. + \vec{E}_0 \cdot \vec{E}_0^* \right]$$

$$= \frac{n \epsilon_0 c}{2} \vec{E}_0 \cdot \vec{E}_0^* \left[ 2 + e^{i[(\vec{k}_2 - \vec{k}_1) \cdot \vec{r} - (\omega_2 - \omega_1)t]} \right. \\ \left. - i[(\vec{k}_2 - \vec{k}_1) \cdot \vec{r} - (\omega_2 - \omega_1)t] \right. \\ \left. + e \right]$$

$$= n \epsilon_0 c \vec{E}_0 \cdot \vec{E}_0^* \left[ 1 + \cos[(\vec{k}_2 - \vec{k}_1) \cdot \vec{r} - (\omega_2 - \omega_1)t] \right]$$

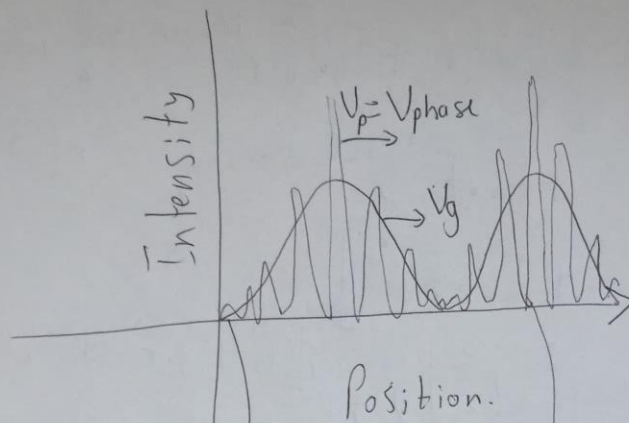
$$= n \epsilon_0 c \vec{E}_0 \cdot \vec{E}_0^* \left[ 1 + \cos(\Delta \vec{k} \cdot \vec{r} - \Delta \omega t) \right]$$

$$\Delta \vec{k} \equiv \vec{k}_2 - \vec{k}_1$$

$$\Delta \omega \equiv \omega_2 - \omega_1$$



(8)



averaged  
over rapid  
oscillations.

keeping  
rapid oscillations.

$$V_p = \frac{\bar{\omega}}{k} \quad ; \quad V_g \equiv \frac{\Delta \omega}{\Delta k} \approx \left. \frac{d\omega}{dk} \right|_{\bar{\omega}}$$

↓  
phase velocity

↓  
group velocity.  
(envelope velocity)

P 7.3 HW

Example 7.1

(9)

$$n_{\text{plasma}}(\omega) = \sqrt{1 - \omega_p^2 / \omega^2} \quad (< 1) \text{ (let } \omega > \omega_p)$$

Find  $V_p, V_g$  for superposition of two plane waves in a plasma.

$$V_p = \frac{\vec{W}}{k} = \frac{W_1 + W_2 / 2}{k_1 + k_2 / 2}$$

$$= \frac{W_1 + W_2}{\frac{n_p(\omega_1) W_1}{c} + \frac{n_p(\omega_2) W_2}{c}}$$

if  $\omega_1 \approx \omega_2 \rightarrow n_p(\omega_1) \approx n_p(\omega_2)$

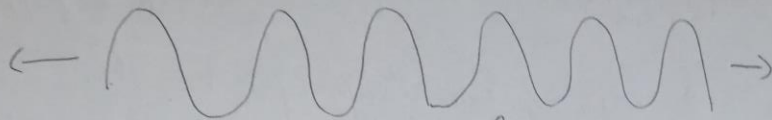
$$V_p = \frac{c}{n_p(\omega)} \quad ?? \quad \boxed{V_p > c}$$

$$V_g = \frac{\Delta \omega}{\Delta k} \approx \frac{d\omega}{dk} = \left[ \frac{dk}{d\omega} \right]^{-1} = \left[ \frac{d}{d\omega} \frac{\omega n_p(\omega)}{c} \right]^{-1}$$

$$= \left[ \frac{d}{d\omega} \frac{\omega \sqrt{1 - \omega_p^2 / \omega^2}}{c} \right]^{-1} = n_p(\omega) c$$

$$\rightarrow \left( 1 - \frac{\omega_p^2}{\omega^2} + \frac{\omega_p^2}{\omega^2} \right) / \sqrt{1 - \frac{\omega_p^2}{\omega^2}}$$

### 7.3 Frequency Spectrum of Light (10)



Plane wave with one frequency  $\omega$  has infinite length and infinite duration.

A waveform that does not repeat, like a pulse is given by:

$$\vec{E}(\vec{r}, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \vec{E}(\vec{r}, \omega) e^{-i\omega t} d\omega$$

spectrum

↓  
gives amplitude  
and phase of each  
plane wave component  
in the waveform.